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Opgave: a) $\frac{d}{dt} \begin{pmatrix} x_t \\ y_t \end{pmatrix} = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} \begin{pmatrix} x_t \\ y_t \end{pmatrix},$

dus $\frac{d}{dt} \vec{v}_t = -i A \cdot \vec{v}_t$ voor $\vec{v}_t = \begin{pmatrix} x_t \\ y_t \end{pmatrix}$

en $-i A = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} \rightsquigarrow A = \begin{pmatrix} 0 & i \\ i & 0 \end{pmatrix}$

$H(\vec{v}) = A \cdot \vec{v}$ is hermitisch omdat $a_{ij} = -\bar{a}_{ji}$,
 $\bar{0} = 0$, $\bar{i} = -i$.

b) Eigenwaarden van $\begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix}$

$\det \begin{pmatrix} \lambda & -1 \\ 1 & \lambda \end{pmatrix} = \lambda^2 + 1 = 0 \quad \lambda = \pm i$

Bij $\lambda = i$: $\vec{v}_i = \begin{pmatrix} 1 \\ i \end{pmatrix}$

Bij $\lambda = -i$: $\vec{v}_{-i} = \begin{pmatrix} 1 \\ -i \end{pmatrix}$

$\begin{pmatrix} x_0 \\ y_0 \end{pmatrix} = c_+ \begin{pmatrix} 1 \\ i \end{pmatrix} + c_- \begin{pmatrix} 1 \\ -i \end{pmatrix}$ geeft

$$\begin{cases} c_+ + c_- = x_0 \\ i(c_+ - c_-) = y_0 \end{cases} \quad \begin{aligned} c_+ &= \frac{1}{2}(x_0 - iy_0) \\ c_- &= \frac{1}{2}(x_0 + iy_0) \end{aligned}$$

Dus $\begin{pmatrix} x_0 \\ y_0 \end{pmatrix} = \frac{1}{2}(x_0 - iy_0) \begin{pmatrix} 1 \\ i \end{pmatrix} + \frac{1}{2}(x_0 + iy_0) \begin{pmatrix} 1 \\ -i \end{pmatrix}$

↓

$\begin{pmatrix} x_t \\ y_t \end{pmatrix} = \frac{1}{2}(x_0 - iy_0) e^{it} \begin{pmatrix} 1 \\ i \end{pmatrix} + \frac{1}{2}(x_0 + iy_0) e^{-it} \begin{pmatrix} 1 \\ -i \end{pmatrix}$

$$= \begin{pmatrix} x_0 \cos t + y_0 \sin t \\ -x_0 \sin t + y_0 \cos t \end{pmatrix}$$

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$$\begin{aligned} c) \quad x_t^2 + y_t^2 &= \left(x_0 \cos(t) + y_0 \sin(t) \right)^2 + \left(y_0 \cos(t) - x_0 \sin(t) \right)^2 \\ &= x_0^2 \cos^2 t + 2x_0 y_0 \cos t \sin t + y_0^2 \sin^2 t \\ &\quad + y_0^2 \cos^2 t - 2x_0 y_0 \cos t \sin t + x_0^2 \sin^2 t \\ &= x_0^2 (\cos^2 t + \sin^2 t) + y_0^2 (\cos^2 t + \sin^2 t) \\ &= x_0^2 + y_0^2 \end{aligned}$$

Opgave: als (x_1, x_2, x_3, x_4) valdeel aan

$$\frac{d}{dt} \begin{pmatrix} x_1 \\ x_2 \\ x_3 \\ x_4 \end{pmatrix} = \begin{pmatrix} 0 & -0,42 & -1,32 & -7,62 \\ 0,42 & 0 & -4,21 & -4,22 \\ 1,32 & 4,21 & 0 & -3,17 \\ 7,62 & 4,22 & 3,17 & 0 \end{pmatrix} \begin{pmatrix} x_1 \\ x_2 \\ x_3 \\ x_4 \end{pmatrix},$$

en $x_1^2 + \dots + x_4^2 = R^2$ op $t=0$, wat 's dat dan
op tyds t ?

Antw: Geen zinnig mens gaan hieraan rekenen.
(of in nder geval geen lui mens.)

De vergelyking is van de vorm $\frac{d}{dt} \vec{v}_t = -iA \cdot \vec{v}_t$,

met $A = \begin{pmatrix} 0 & -i \cdot 0,42i & -1,52i & \cdots \\ 0,42i & 0 & -4,21i & \cdots \\ 1,32i & 4,21i & 0 & \cdots \\ \vdots & \vdots & \vdots & \ddots \end{pmatrix}$

A is hermitosech, $a_{ij} = \bar{a}_{ij}$.

Dus $\langle \vec{v}_t, \vec{v}_0 \rangle = \langle \vec{v}_0, \vec{v}_0 \rangle$.

Maar x_1, \dots, x_4 is reëel, dus x_1

$$\langle \vec{v}_{t=0}, \vec{v}_{t=0} \rangle = |x_1(0)|^2 + \dots + |x_4(0)|^2 = x_1^2(0) + \dots + x_4^2(0) = R^2$$

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Omdat $\frac{dt}{dt} x_i$ reeel is voor alle t is ook
 $x_i(t)$ reeel, en $x_i^2(t)$

$$\langle \vec{v}_t, \vec{v}_t \rangle = |x_1(t)|^2 + \dots + |x_n(t)|^2 = x_1^2(t) + \dots + x_n^2(t).$$

Omdat $\langle v(t=0), \vec{v}(t=0) \rangle = \langle \vec{v}(t), \vec{v}(t) \rangle$,

is dus $x_1^2(t) + \dots + x_n^2(t) = R^2$.

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Opg: Voor een vast getal c_g behalen we

$$\frac{d}{dt} \Psi_t = -c_g \frac{d}{d\phi} \Psi_t$$

a) $\frac{d}{dt} \Psi_t = L \Psi_t$ met $L = -c_g \frac{d}{d\phi}$

$$L \Psi_t = -c_g \frac{d}{d\phi} e^{i c_g t \phi} = -i c_g k e^{i c_g t \phi}$$

dus $\lambda_k = -i c_g k$ is de eigenwaarde, $w_k = \underline{c_g k}$

c) $\Psi_k(\phi, t) = e^{\lambda_k t} \Psi_k(\phi) = e^{-i c_g k t} \cdot e^{i c_g k \phi}$

Met $\Psi_0(\phi)$

d) $\Psi_0 = \lim_{N \rightarrow \infty} \sum_{k=-N}^N c_k e^{i c_g k \phi}$

Dan $\Psi_t(\phi) = \lim_{N \rightarrow \infty} c_k e^{-i c_g k t} e^{i c_g k \phi}$,

dus $c_k(t) = \underline{c_k e^{-i c_g k t}}$.

e) $\Psi_t(\phi) = \lim_{N \rightarrow \infty} \sum_{k=-N}^N c_k \cdot e^{-i c_g k t} \cdot e^{i c_g k \phi}$

$$= \lim_{N \rightarrow \infty} \sum_{k=-N}^N c_k e^{i \lambda_k (\phi - c_g t)}$$

$$= \Psi_0(\phi - c_g t)$$

b) \Rightarrow

$$V = \omega_r / \ell = c_g$$

f) $H = -c_g i \frac{d}{d\phi}$ is een reëel veelvoud van de momentumsoperator $P = -i \hbar \frac{d}{d\phi}$

Inderdaad is $\|\Psi_t\|^2 = \int_{-\pi}^{\pi} |\Psi(\phi - c_g t)|^2 d\phi = \int_{-\pi}^{\pi} |\Psi(\phi)|^2 d\phi = \|\Psi_0\|^2$.

ψ_h

Opgave : Laat zien dat $H = -\frac{t^2}{2m}(\partial_x^2 + \partial_y^2 + \partial_z^2) + V$ hermitisch is.

We kunnen zien dat $P_x = -it\partial_x$ hermitisch is.

$$\int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} (-it\partial_x \psi_1) \psi_2 dx dy dz =$$

$$it \cdot \int_{-\infty}^{\infty} dy \int_{-\infty}^{\infty} dz \left(\int_{-\infty}^{\infty} \partial_x \bar{\psi}_1 \cdot \psi_2 dx \right)$$

$$\hookrightarrow = \left[[\bar{\psi}_1 \psi_2] \right]_{x=-\infty}^{x=\infty} - \int_{-\infty}^{\infty} \bar{\psi}_1 \partial_x \psi_2(x, y, z) dz$$

$$\hookrightarrow = 0$$

$$= -it \int_{-\infty}^{\infty} dy \int_{-\infty}^{\infty} dz \int_{-\infty}^{\infty} (\bar{\psi}_1 \partial_x \psi_2) dx =$$

$$\int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \bar{\psi}_1(x, y, z) \partial_x \psi_2 \cdot -it \partial_x \psi_2(x, y, z) dx dy dz.$$

$$\text{Dus } \langle P_x \psi_1, \psi_2 \rangle = \langle \psi_1, P_x \psi_2 \rangle.$$

Net zo zijn P_y en P_z hermitisch.

Dus P_x^2, P_y^2, P_z^2 zijn hermitisch.

Dus het reële veelvoud van hun som,

$$\frac{t^2}{2m} (\partial_x^2 + \partial_y^2 + \partial_z^2)$$

is hermitisch.

$$\begin{aligned} \text{Omdat } V \text{ reeel is is } \langle \psi_1, V \psi_2 \rangle &= \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \bar{\psi}_1 V \psi_2 dx dy dz \\ &= \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \bar{\psi}_1 \psi_2 dx dy dz = \langle V \psi_1, \psi_2 \rangle \end{aligned}$$

Dus V is ook hermitisch, en $H = \frac{t^2}{2m} (\partial_x^2 + \partial_y^2 + \partial_z^2) + V$ is hermitisch.